I wonder if what is done here is totally correct.

6.2 INTERFERENCE OF TWO WAVES

Imagine that we have two linearly polarized plane waves of the same wavelength, given by

$$\mathbf{E}_{1}(\mathbf{r}, t) = \mathbf{E}_{01} \cos \left(\mathbf{k}_{1} \cdot \mathbf{r} - \omega t + \varepsilon_{1}\right)$$
$$\mathbf{E}_{2}(\mathbf{r}, t) = \mathbf{E}_{02} \cos \left(\mathbf{k}_{2} \cdot \mathbf{r} - \omega t + \varepsilon_{2}\right)$$

which overlap at point P as in Fig. 6-1. Here, of course, \mathbf{k}_1 , \mathbf{k}_2 , ω , ε_1 and ε_2 are all constant. These waves may arise, for example, from two, very distant, point sources. The resultant field is simply

 $\mathbf{E} = \mathbf{E}_1 + \mathbf{E}_2$

But since at optical frequencies the fields oscillate at in excess of 10^{14} Hz the irradiance becomes the quantity of practical concern, it being directly measurable. Neglecting a constant factor, we can write the irradiance as just the time average of the total field:

 $I = \langle \mathbf{E}^2 \rangle$

where $\mathbf{E}^2 \equiv \mathbf{E} \cdot \mathbf{E}$. Accordingly,

$$\mathbf{E}^2 = (\mathbf{E}_1 + \mathbf{E}_2) \cdot (\mathbf{E}_1 + \mathbf{E}_2) = \mathbf{E}_1^2 + \mathbf{E}_2 + 2\mathbf{E}_1 \cdot \mathbf{E}_2$$

Taking the average we obtain

$$I = I_1 + I_2 + I_{12}$$

where $I_1 = \langle \mathbf{E}_1^2 \rangle$, $I_2 = \langle \mathbf{E}_2^2 \rangle$ and $I_{12} = 2 \langle \mathbf{E}_1 \cdot \mathbf{E}_2 \rangle$, the last being known as the *interference term*. It is by virtue of this last term that I differs from a simple sum of the irradiances of the component waves, i.e. $I_1 + I_2$.

The time average of some function f(t) is given by

$$\langle f(t) \rangle = \frac{1}{T} \int_{t}^{t+T} f(t') dt'$$

Here the detection interval T is much greater than the period of oscillation of the waves, τ . Thus if we carried out the indicated calculation for the above plane waves (see Problem 2.13), the interference term would become

$$I_{12} = 2\langle \mathbf{E}_1 \cdot \mathbf{E}_2 \rangle = \mathbf{E}_{01} \cdot \mathbf{E}_{02} \cos \delta$$

where the phase difference δ is given by

$$\delta = (\mathbf{k}_1 \cdot \mathbf{r}) - (\mathbf{k}_2 \cdot \mathbf{r}) + \varepsilon_1 - \varepsilon_2$$

All of this means that as one moves from point to point in space, r varies, as does δ , and therefore I_{12} and I both vary as well.

I try to explain my doubts. We define intensity as the power transferred per unit area perpendicular to the direction of propagation of the energy. If we have a plane wave this doesn't give problems, but here we have two waves and it is not clear which is the plane. So I think we should define intensity as a vector \mathbf{S} such that $\mathbf{S} \cdot d\mathbf{a}dt = dE$, where dE is the energy that cross $d\mathbf{a}$ during $t \to t + dt$. Maxwell theory tell us that a such vector exists, and its value is $\mathbf{S} = \frac{\mathbf{E} \times \mathbf{B}}{\mu_0}$ and that if we have a plane wave $\mathbf{S} = c\epsilon_0 E^2 \hat{\mathbf{n}}$ where $\hat{\mathbf{n}} = \hat{\mathbf{E}} \times \hat{\mathbf{B}}$ (I use hat for versors and I wrote instantaneous values, I have not yet averaged over rapidly oscillating waves). But I can't understand how the book treat the case of two waves.

Let's try to work the problem: in \mathbf{r} at time t the fields are $\mathbf{E} = \mathbf{E}_1 + \mathbf{E}_2$ and $\mathbf{B} = \mathbf{B}_1 + \mathbf{B}_2$ (I use subscript to denote the wave). Calculating \mathbf{S} we get

$$\mathbf{S} = \mathbf{S}_1 + \mathbf{S}_2 + \frac{1}{\mu_0} (\mathbf{E}_1 \times \mathbf{B}_2 + \mathbf{E}_2 \times \mathbf{B}_1)$$

But if $\hat{\mathbf{n}}$ is the direction of an electromagnetic wave $\mathbf{B} = \frac{\hat{\mathbf{n}} \times \mathbf{E}}{c}$ so

$$\mathbf{S} = \mathbf{S}_1 + \mathbf{S}_2 + \epsilon_0 c [\mathbf{E}_1 \times (\mathbf{\hat{n}}_2 \times \mathbf{E}_2) + \mathbf{E}_2 \times (\mathbf{\hat{n}}_1 \times \mathbf{E}_1)]$$

This can be rearranged with the triple vector product identity, but the result it doesn't look nicer or useful, so I leave it in that way. If the two waves are harmonic, $\mathbf{E}_1 = \mathbf{E}_{01} \cos(\mathbf{k}_1 \cdot \mathbf{r} - \omega t + \epsilon_1)$ and $\mathbf{E}_2 = \mathbf{E}_{02} \cos(\mathbf{k}_2 \cdot \mathbf{r} - \omega t + \epsilon_2)$, and we get

$$\mathbf{S} = \mathbf{S}_1 + \mathbf{S}_2 + \mathbf{X}\cos(\mathbf{k}_1 \cdot \mathbf{r} - \omega t + \epsilon_1)\cos(\mathbf{k}_2 \cdot \mathbf{r} - \omega t + \epsilon_2)$$

where I used the *constant* vector

$$\mathbf{X} = \epsilon_0 c [\mathbf{E}_{01} \times (\mathbf{\hat{n}}_2 \times \mathbf{E}_{02}) + \mathbf{E}_{02} \times (\mathbf{\hat{n}}_1 \times \mathbf{E}_{01})]$$

So

$$<\mathbf{S}>=<\mathbf{S}_1>+<\mathbf{S}_2>+\mathbf{X}<\cos(\mathbf{k}_1\cdot\mathbf{r}-\omega t+\epsilon_1)\cos(\mathbf{k}_2\cdot\mathbf{r}-\omega t+\epsilon_2)>$$

Exploiting $\cos(a-b) = \cos a \cos b + \sin a \sin b$ we have that $\langle \cos [(\mathbf{k_1} \cdot \mathbf{r} + \epsilon_1) - \omega t] \cos [(\mathbf{k_2} \cdot \mathbf{r} + \epsilon_2) - \omega t] \rangle$ is

$$<\left[\cos(\mathbf{k_1}\cdot\mathbf{r}+\epsilon_1)\cos(\omega t)+\sin(\mathbf{k_1}\cdot\mathbf{r}+\epsilon_1)\sin(\omega t)\right]\cdot\left[\cos(\mathbf{k_2}\cdot\mathbf{r}+\epsilon_2)\cos(\omega t)+\sin(\mathbf{k_2}\cdot\mathbf{r}+\epsilon_2)\sin(\omega t)\right]>(1)$$

But if the wave oscillates rapidly we can exploit $\langle \cos^2(\omega t) \rangle = \langle \sin^2(\omega t) \rangle = \frac{1}{2}$ and $\langle \cos(\omega t) \sin(\omega t) \rangle = 0$, so doing products and observing again $\cos(a - b) = \cos a \cos b + \sin a \sin b$ we conclude that

$$<\mathbf{S}> = <\mathbf{S}_1> + <\mathbf{S}_2> + \mathbf{X}\cos(\mathbf{k_1}\cdot\mathbf{r} + \epsilon_1 - \mathbf{k_2}\cdot\mathbf{r} - \epsilon_2)$$

This recalls the book equation, but it is not the same because I use vectors and because the I_{12} term is very different (I used the X vector, ugly, but isn't it necessary?). So I have two question:

- Is it correct in this two waves problem to treat intensity as a scalar, as book does? (the books has a general approach, it doesn't make hypothesis like $\hat{\mathbf{n}}_1 \approx \hat{\mathbf{n}}_2$, on the contrary a figure explicitly show that waves moves in very different directions)
- Why my equation is so different from the book's one, where did I wrong? Is it possible simply find intensity averaging $|\mathbf{E}|^2$ as we do when we have a single wave?

Edit

The key is that Poynting says that the only important thing to find intensity vector in a given point at a given time are total **E** and **B**, but if we have many waves overlapping the total **E** and **B** are *in some way* independent. I mean that they can have any relative direction and the ratio of magnitude is not c. How can we so conoclude that, in a generic overlapping of two waves, the intensity in a give place **r** and time t is proportional to the square of local electric field? This can be demonstrated for one wave (by exploiting $\mathbf{E} \perp \mathbf{B}$ and E = Bc) but not in a general overlapping of waves, as it is assumed in the book proof.

By the way, if the two waves moves approximately in the same direction $(\hat{\mathbf{n}}_1 \approx \hat{\mathbf{n}}_2 = \hat{\mathbf{n}})$, the book reasoning gives approximately the correct result (using triple vector product identity with my vector \mathbf{X} and the fact that $\mathbf{E}_{01}, \mathbf{E}_{02} \perp \hat{\mathbf{n}}$), but this hypothesis is not assumed, the approach is general, so it looks an error. I wrong?